Superconductivity - Overview

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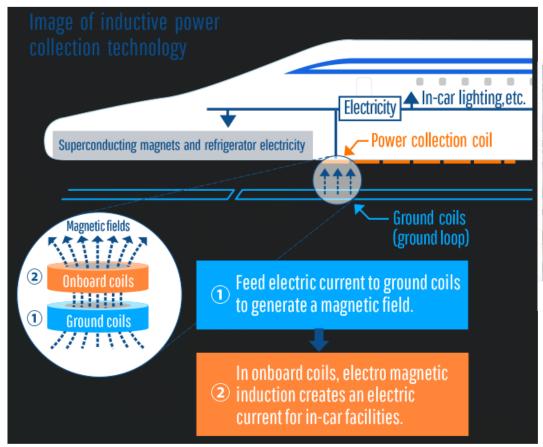
Importance of superconductivity: Qubits and modern quantum computers are made of superconductors

IBM Q





Superconducting-magnet levitation train





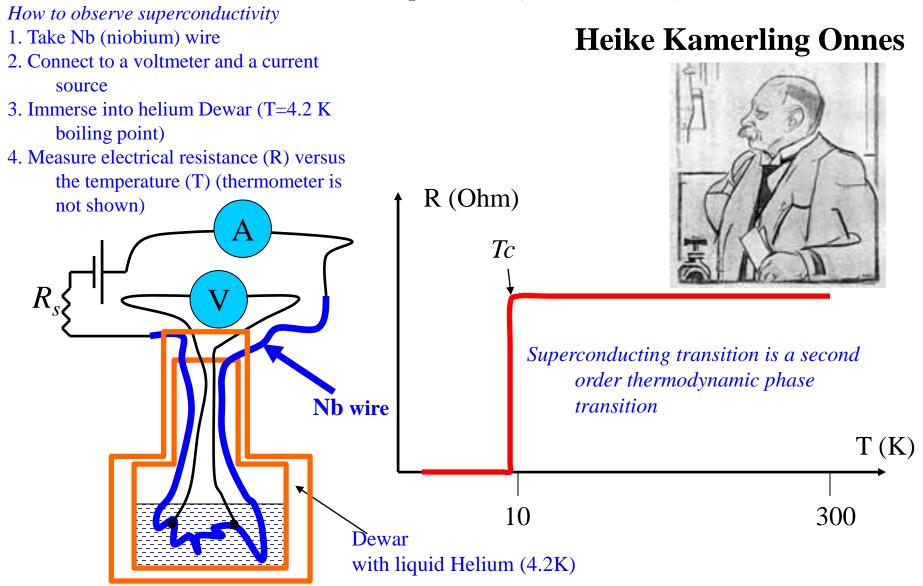
The <u>Lo Series</u>, a prototype vehicle based on SCMaglev technology, holds the record for fastest crewed rail vehicle with a record speed of 603 km/h (375 mph).

Time Urbana-Chicago: only ~25 min

The SCMaglev system uses an electrodynamic suspension (EDS) system. The train's bogies have superconducting magnets installed, and the guideways contain two sets of metal coils.

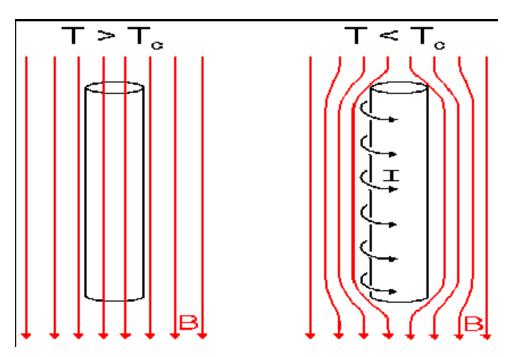
How to measure superconducting transitions

Electrical resistance of some metals drops to zero below a certain temperature which is called "critical temperature" (H. K. O. 1911)





Meissner effect – the key signature of superconductivity

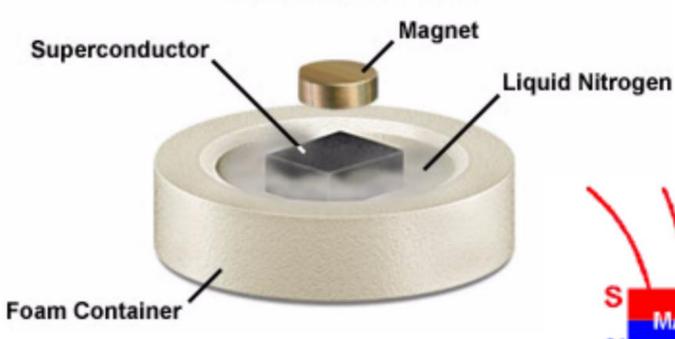


Theory of superconductivity: "BCS" – due to Bardeen, Cooper and Schrieffer

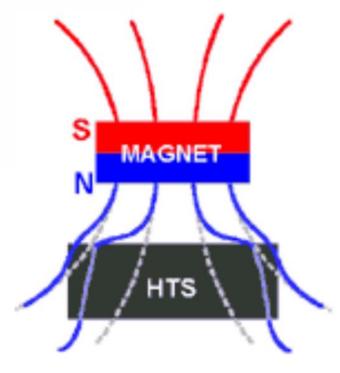
| Formula | т _с (K) | H _C (T) | Туре | BCS |
|-----------|-----------------------|-----------------------|------|-----|
| Elements | | | | |
| Al | 1.20 | 0.01 | I | yes |
| Cd | 0.52 | 0.0028 | I | yes |
| Diamond:B | 11.4 | 4 | II | yes |
| Ga | 1.083 | 0.0058 | I | yes |
| Hf | 0.165 | | I | yes |
| α-Hg | 4.15 | 0.04 | I | yes |
| β-Нд | 3.95 | 0.04 | I | yes |
| In | 3.4 | 0.03 | I | yes |
| Ir | 0.14 | 0.0016 ^[7] | I | yes |
| α-La | 4.9 | | I | yes |
| β-La | 6.3 | | I | yes |
| Мо | 0.92 | 0.0096 | I | yes |
| Nb | 9.26 | 0.82 | II | yes |
| Os | 0.65 | 0.007 | I | yes |

Interesting phenomenon: Magnetic levitation

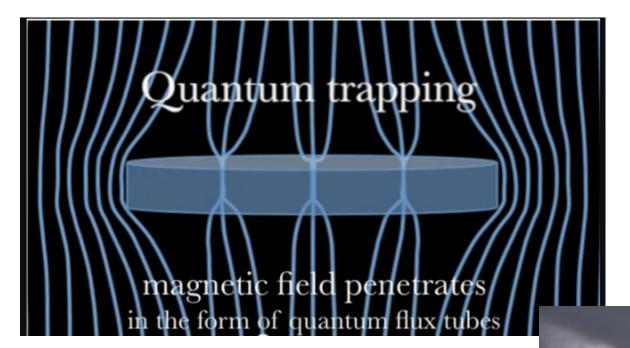
The Meissner Effect



Levitation is the process by which an object is held aloft, without mechanical support, in a stable position.



Magnetic field effect: Superconducting vortices



In superconductivity, a

fluxon (also called an

Abrikosov vortex or

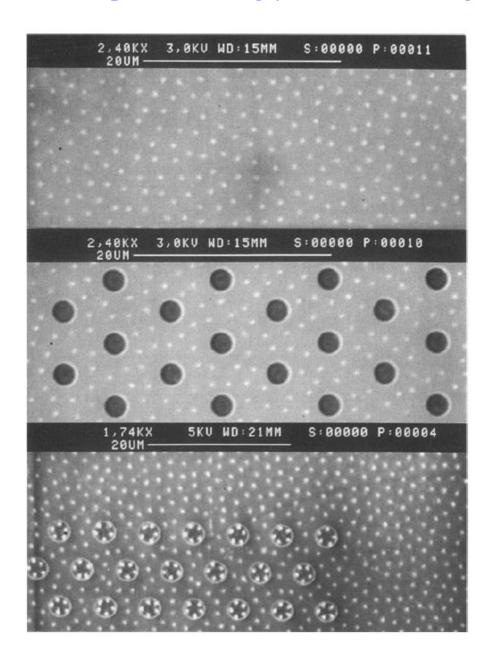
quantum vortex) is a

vortex of supercurrent in

a type-II superconductors

https://blog.tmcnet.com/blog/tomkeating/technology-and-science/quantumlevitation-back-to-the-futurehoverboard.asp

Vortices in superconducting films with "through" and "blind" holes ("antidots")

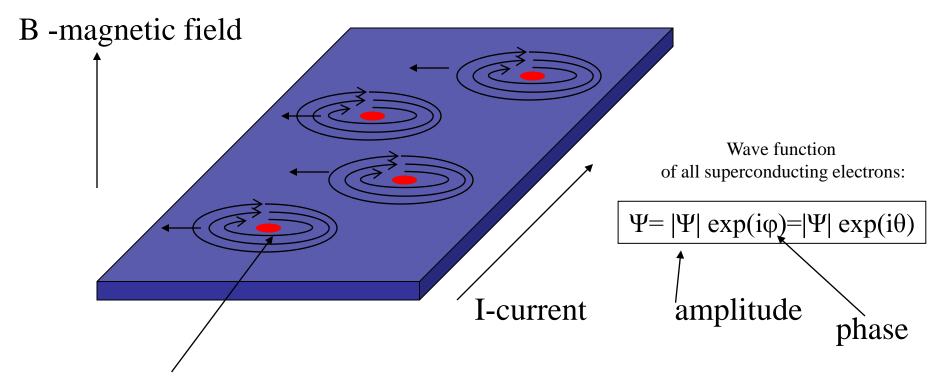


A. Bezryadin and B. Pannetier
"Role of Edge Superconducting States in
Trapping of Multi-Quanta Vortices by
Microholes. Application of the Bitter
Decoration Technique",
J. Of Low Temp. Phys., V.102, p.73 (1996).

Vortices are quantized tubes carrying magnetic field into superconductor

Magnetic field creates vortices--

Vortices cause dissipation (i.e. a non-zero electrical resistance), if they move



Vortex core (red): normal, not superconducting; diameter $\xi\sim10$ nm The current is extended to a scale λ , which is larger than ξ in type II superconductors (such as thin films of any material)

Reminder: single electron in empty space

Wave function:
$$\Psi = |\Psi| \exp(ikx) = |\Psi| [\cos(kx) + i \sin(kx)]$$

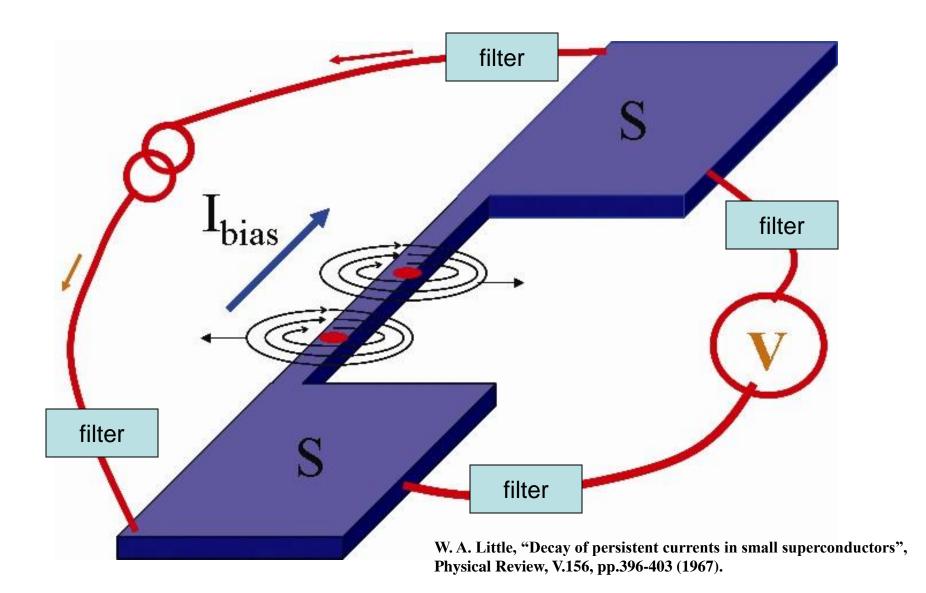
Wave number: $k=2\pi/\lambda$
 $i*i=-1$

General form: $\Psi = |\Psi| \exp(-i\varphi)$ In this example of a plane wave, the phase is: $\varphi = kx$

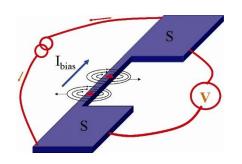
DC transport measurement schematic to detect passing vortices

Bleu: superconducting film and wire

Red: Phase slip events or crossing vortices



How to use voltage to determine the rate of phase slips?



Key principle: every time a vortex crosses the wire the phase difference changes by 2π .

Phase evolution equation: $d\phi/dt = 2eV/\hbar$

Simplified derivation:

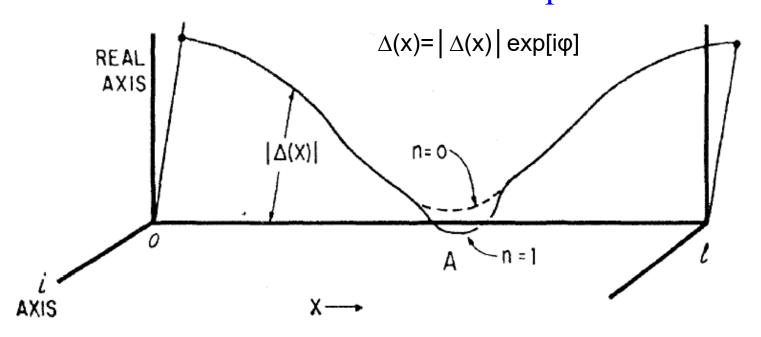
1. Time-dependent Schrödinger equation with fixed energy:

 $i\hbar(d\Psi/dt)=E\Psi$

- 2. The solution is: $\Psi = \exp(-iEt/\hbar)$ (here E is the energy)
- 3. The phase of the wavefunction is $\varphi = Et/\hbar$
- 4. The energy is defined by the electric potential (voltage), V as follows: E=2eV. Note that the effective charge of superconducting electrons is 2e, where "e" is the charge of one electron. Such superconducting electron pairs are called Cooper pairs.

Thus, the resulting equation is: $d\phi/dt = 2eV/\hbar_{Gor'kov, L.P. (1958) Sov. Phys. JETP, 7,505.}$

Thin superconducting wire have some nonzero electrical resistance due to Little's Phase Slips



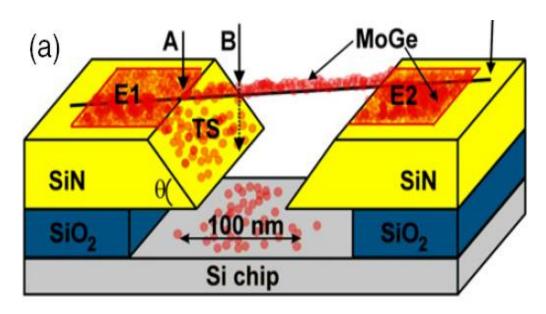
W. A. Little, "Decay of persistent currents in small superconductors", Physical Review, V.156, pp.396-403 (1967).

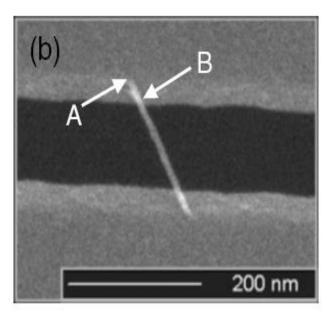
Two types of phase slips (PS) can occur:

- 1. The usual, thermally activated PS (TAPS)
- 2. Quantum phase slip (QPS)

Fabrication of nanowires

Method of Molecular Templating





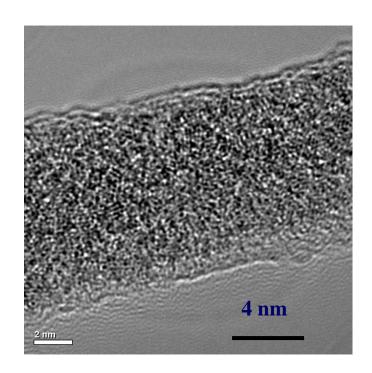
Si/SiO₂/SiN substrate with undercut

 $\sim 0.5 \text{ mm Si wafer}$ 500 nm SiO_2 60 nm SiNWidth of the trenches ~ 50 - 500 nm

HF wet etch for ~10 seconds to form undercut



Sample Fabrication



100 nm nanowire nanotube

SiO₂

SiO₂

TEM image of a wire shows amorphous wire morphology.

Nominal MoGe thickness = 3 nm

Schematic picture of the pattern Nanowire + Film Electrodes used in transport measurements



Dichotomy in nanowires: Evidence for superconductorinsulator transition (SIT)

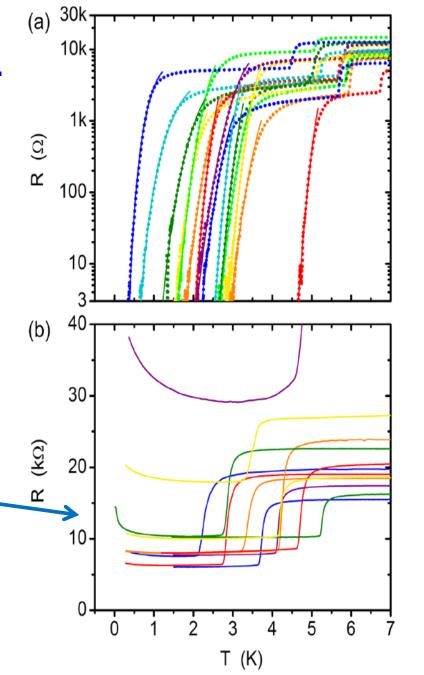
R=V/I $I\sim3$ nA

The difference between samples is the amount of the deposited Mo79Ge21.

Thin wires become insulating if their normal resistance is lager than resistance quantum h/4/e/e=6.5KOhm

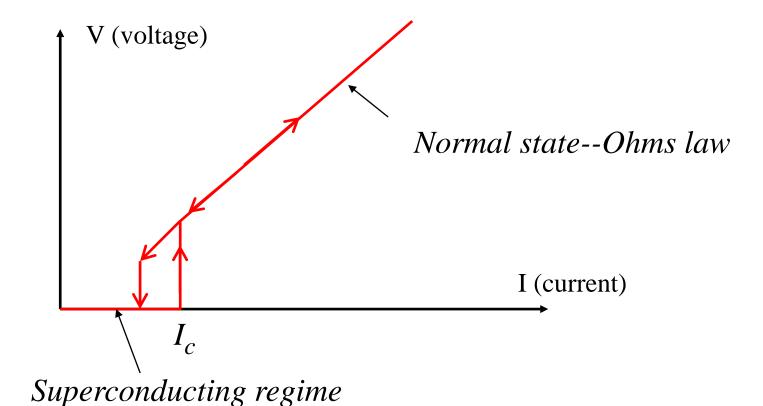
The insulating behavior is due to proliferation of quantum phase slips

Bollinger, Dinsmore, Rogachev, Bezryadin, Phys. Rev. Lett. **101**, 227003 (2008)



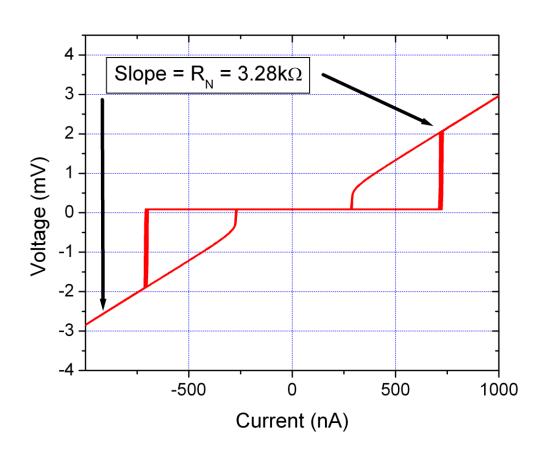
Expected voltage-current curve

Electrical resistance is zero only if current is not too strong



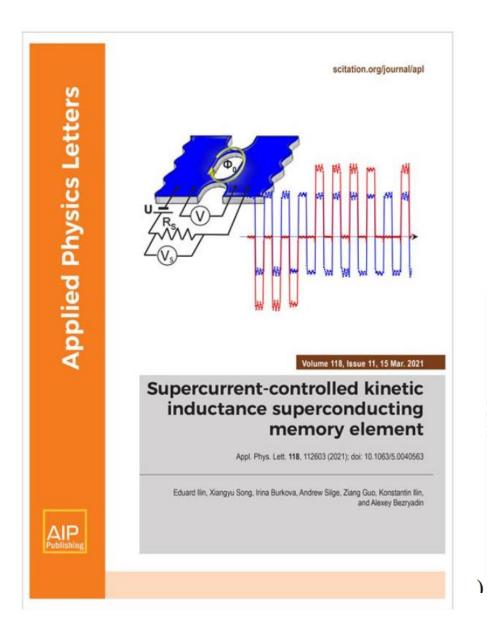
Experimental voltage-current curve.

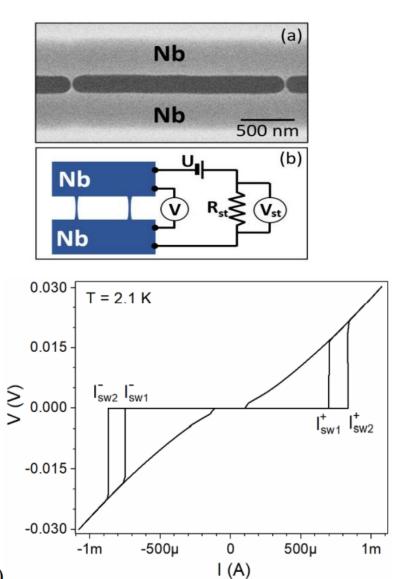
Fluctuations of the switching current are due to Little's phase slips





Superconducting nanowire memory





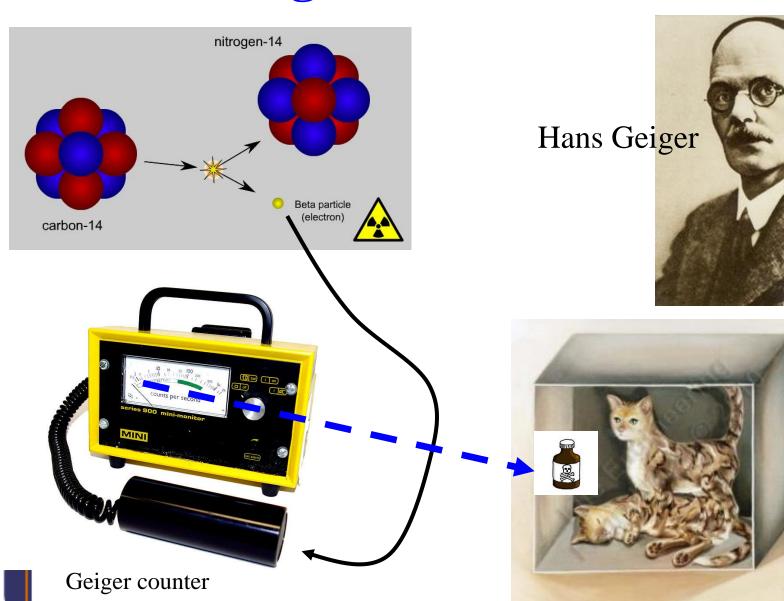
Schrödinger cat – the ultimate macroscopic quantum phenomenon

E. Schrödinger, Naturwiss. 23 (1935), 807.





$Schr\"{o}dinger\ cat\ -\ thought\ experiment$





Suppose Ψ_1 is a valid solution of the Schrödinger equation:

$$i\hbar \frac{\partial \psi_1}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi_1}{\partial x^2} + U(x)\psi_1$$

And suppose that Ψ_2 is another valid solution of the Schrödinger equation:

$$i\hbar \frac{\partial \psi_2}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi_2}{\partial x^2} + U(x)\psi_2$$

Then $(\Psi_1 + \Psi_2)/\sqrt{2}$ is also a valid solution, because:

$$i\hbar \frac{\partial (\psi_1 + \psi_2)}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 (\psi_1 + \psi_2)}{\partial x^2} + U(x)(\psi_1 + \psi_2)$$

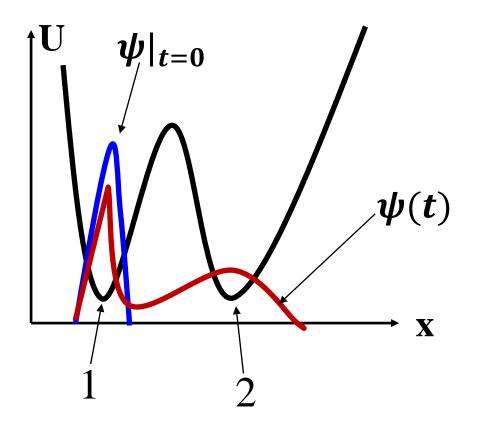
The state $(\Psi_1 + \Psi_2)/\sqrt{2}$ is a new combined state which is called "quantum superposition" of state (1) and (2)



George Gamow

(He also helped to developed Big Bang theory)

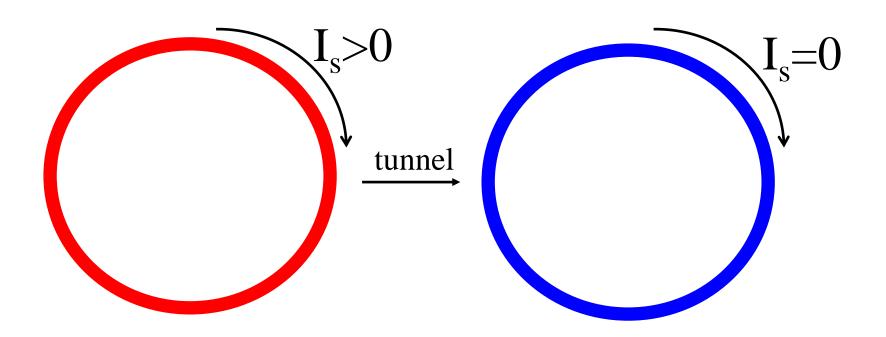
Quantum tunneling



Quantum tunneling is possible since quantum superpositions of states are possible.



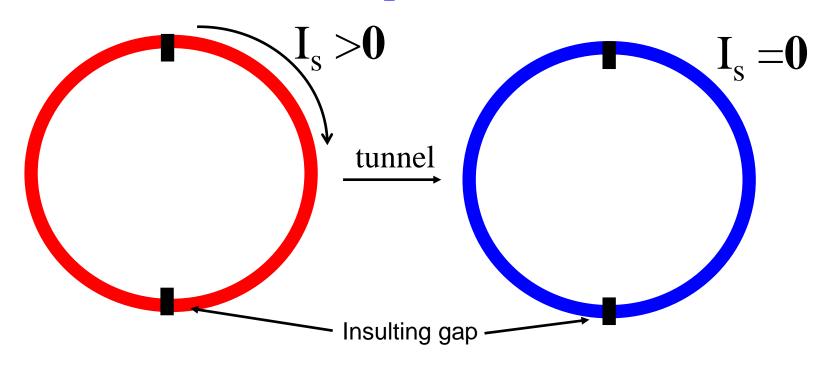
What sort of tunneling we will consider?



- -Red color represents some strong current in the superconducting wire loop
- -Blue color represents zero current in the loop



Previous results relate loops with insulating interruptions (SQUIDs)



- -Red color represents some strong current in the superconducting loop
- -Blue color represents zero current in the superconducting loop



Leggett's prediction for macroscopic quantum tunneling (MQT) in SQUIDs

80

Supplement of the Progress of Theoretical Physics, No. 69, 1980

Macroscopic Quantum Systems and the Quantum Theory of Measurement

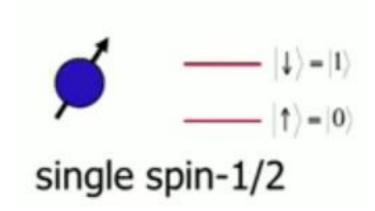
A. J. LEGGETT

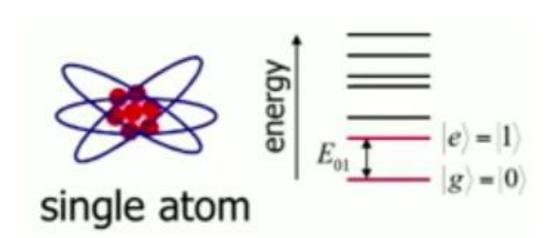
School of Mathematical and Physical Sciences University of Sussex, Brighton BN1 9QH

(Received August 27, 1980)

It is this property which makes a SQUID the most promising candidate to date for observing macroscopic quantum tunnelling; if it should ever become possible to observe macroscopic quantum coherence, the low entropy and consequent lack of dissipation will be absolutely essential.²¹⁾

Types of Qubit



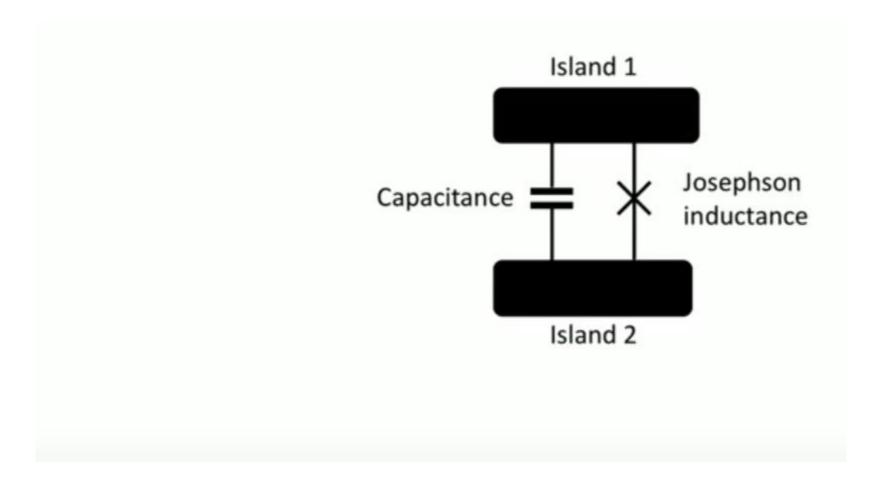


Quantum state: $|\psi\rangle = A^*|0\rangle + B^*1\rangle$

$$A^2 + B^2 = 1$$

A and B are complex numbers

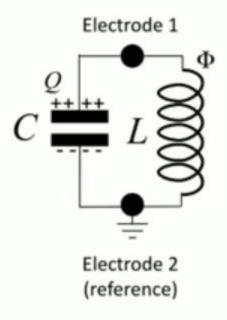
Transmon Qubit



Theory of transmons: J. Koch et al., Phys. Rev. A 76, 042319 (2007).

Quantization of electrical circuits

The quantized *LC* oscillator



Hamiltonian:

$$\hat{H}_{LC} = \frac{\hat{Q}^2}{2C} + \frac{\hat{\Phi}^2}{2L}$$
Capacitive term Inductive term

Canonically conjugate variables:

$$\hat{\Phi}$$
 = Flux through the inductor.

$$\hat{Q}$$
 = Charge on capacitor plate.

$$\left[\hat{\Phi},\hat{Q}\right] = i\hbar$$

M. Devoret, Les Houches Session LXIII (1995)

Discrete energy spectrum of the LC-circuit

Correspondence with simple harmonic oscillator

$$\hat{H}_{\mathrm{LC}} = \frac{\hat{\Phi}^2}{2L} + \frac{\hat{Q}^2}{2C}$$

$$\left[\hat{\Phi}, \hat{Q}\right] = i\hbar$$

$$\left[\hat{X}, \hat{P}\right] = i\hbar$$

$$\left[\hat{X}, \hat{P}\right] = i\hbar$$

$$\left[\hat{X}, \hat{P}\right] = i\hbar$$
 Correspondence:
$$\hat{Q} \leftrightarrow \hat{P}$$

$$C \leftrightarrow m$$

$$\omega = \frac{1}{\sqrt{LC}} \leftrightarrow \sqrt{\frac{k}{m}}$$

$$\omega = \frac{1}{\sqrt{LC}} \leftrightarrow \sqrt{\frac{k}{m}}$$

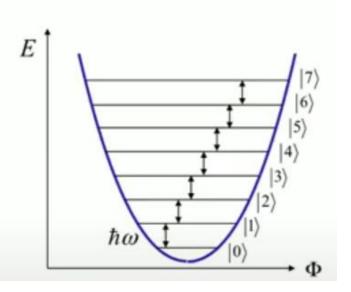
Solve using ladder operators:

$$\hat{a} = \left(\frac{\hat{Q}}{Q_{\eta\ell}} - i\frac{\hat{\Phi}}{\Phi_{\eta\ell}}\right) \qquad \Phi_{\eta\ell} = \sqrt{2\hbar Z}$$

$$\hat{a}^{\dagger} = \left(\frac{\hat{Q}}{Q_{\eta\ell}} + i\frac{\hat{\Phi}}{\Phi_{\eta\ell}}\right) \qquad Z = \omega L = \frac{1}{\omega C} = \sqrt{\frac{L}{C}}$$

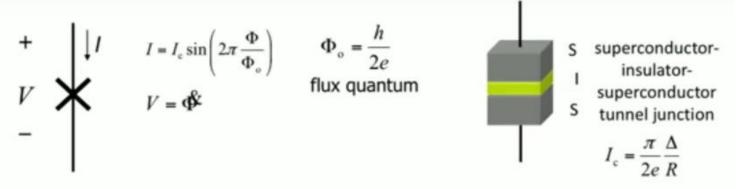
$$\hat{H}_{LC} = \hbar \omega \left(\hat{a}^{\dagger} \hat{a} + \frac{1}{2}\right) \qquad \left[\hat{a}_{r}, \hat{a}_{r}^{\dagger}\right] = 1$$

M. Devoret, Les Houches Session LXIII (1995)



Non-harmonicity is the key factor

The Josephson junction

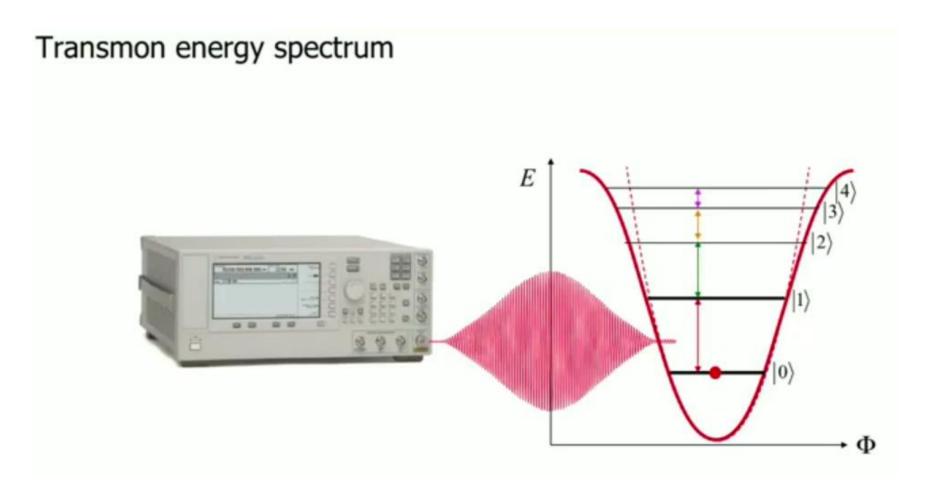


$$E_{\rm stored} = E_{\rm J} \left(1 - \cos \left(2\pi \frac{\Phi}{\Phi_o} \right) \right) \qquad E_{\rm J} = \frac{I_{\rm c} \Phi_o}{2\pi} \quad \text{Josephson}$$
 Energy
$$E_{\rm J} = \frac{I_{\rm c} \Phi_o}{2\pi} \quad \text{Josephson potential}$$
 For small flux:
$$E_{\rm J} = \frac{\Phi^2}{2L_{\rm J}} + {\rm const}$$

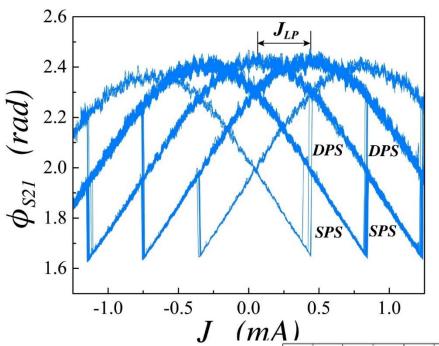
$$\Phi_o$$
 Josephson inductance
$$L_{\rm J} = \frac{\Phi_o}{2\pi I_{\rm c}}$$

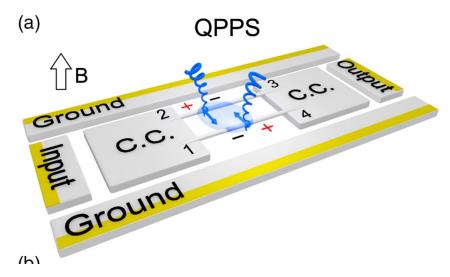
M. Devoret, Les Houches Session LXIII (1995)

Non-harmonicity is the key factor



Superconducting nanowire memory: microwave readout



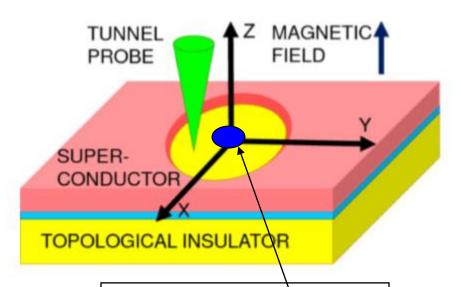


A. Belkin et al, PRX 5, 021023 (2015)

Andrey Belkin,a Matthew Brenner, Thomas Aref, Jaseung Ku, and Alexey Bezryadin, PPLIED PHYSICS LETTERS **98**, 242504 (2011)



Majorana modes in a vortex



Theory: Vortex in the nano-hole contains
Majorana states (gap is about Tc)

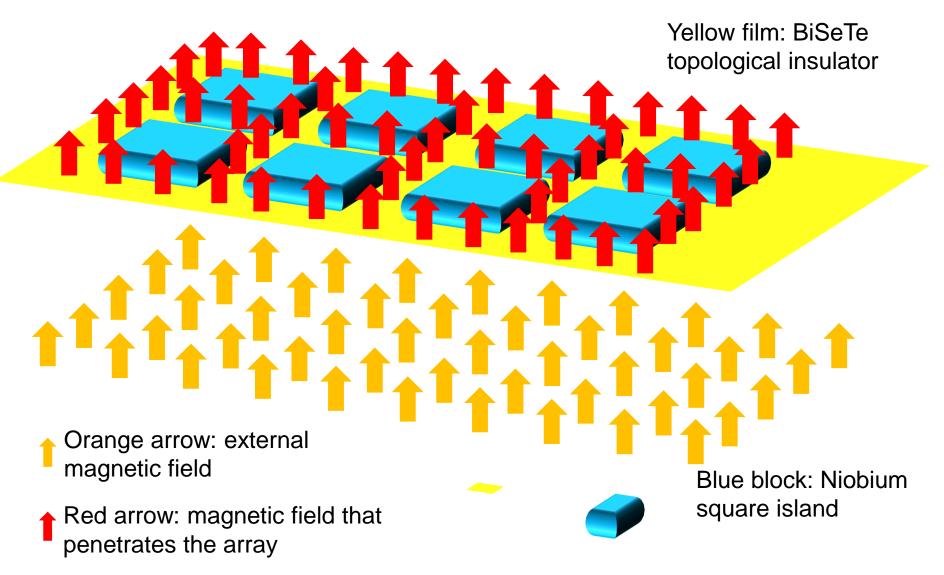
R.S. Akzyanov, A.V. Rozhkov, A.L.Rakhmanov, and F. Nori, PRB 89, 1085409 (2014) 3,94KX 5KU WD:29MM \$:00000 P:00011

A. Bezryadin, Yu. Ovchinnikov, B. Pannetier, PRB 53, 8553 (1996)



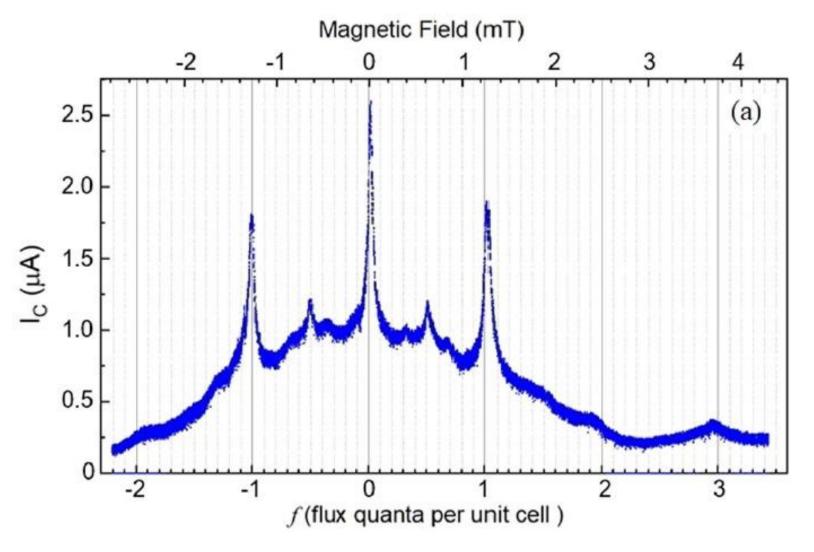
PHYSICAL REVIEW B 84, 075141 (2011)

Superconducting array based on topological insulator- BST



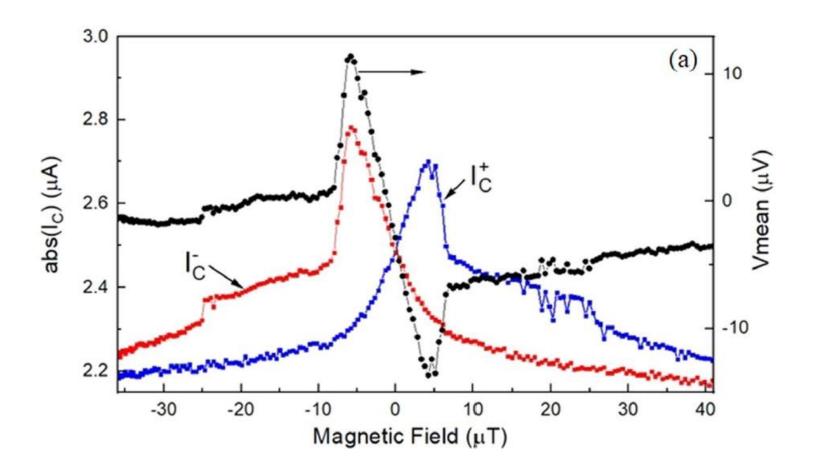
Xiangyu Song, Soorya Suresh Babu, Yang Bai, Dmitry S. Golubev, Irina Burkova, Alexander Romanov, Eduard Ilin, James N. Eckstein & Alexey Bezryadin

Diffraction grating analogy with superconducting arrays



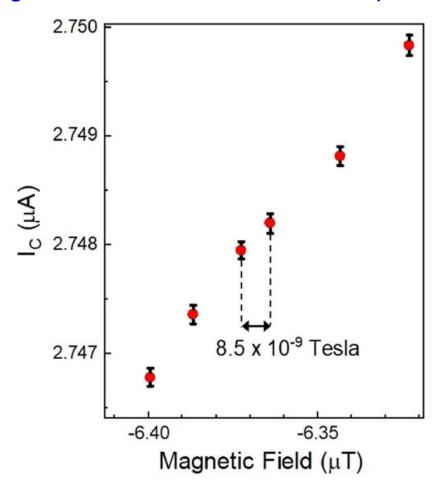
Xiangyu Song, Soorya Suresh Babu, Yang Bai, Dmitry S. Golubev, Irina Burkova, Alexander Romanov, Eduard Ilin, James N. Eckstein & Alexey Bezryadin

Diode effect with superconducting arrays



Xiangyu Song, Soorya Suresh Babu, Yang Bai, Dmitry S. Golubev, Irina Burkova, Alexander Romanov, Eduard Ilin, James N. Eckstein & Alexey Bezryadin

Absolute magnetic field sensor based on superconducting array



Xiangyu Song, Soorya Suresh Babu, Yang Bai, Dmitry S. Golubev, Irina Burkova, Alexander Romanov, Eduard Ilin, James N. Eckstein & Alexey Bezryadin

Conclusions

- Superconductivity is related to fundamental quantum phenomena. We have revied some of them. They will be discussed in more details in the future lectures.
- Superconductors have been used to create strong and stable magnetic fields, in levitating trains for example.
- Superconducting quantum interference devices enabled researches to measure very small magnetic fields, such as those produced by human brain.
- Superconductors are used to build qubits, which are the building blocks of quantum computers.

